ON THE VARIABILITY OF HEMISPHERIC SCALE ENERGY PARAMETERS

by

J. P. McGuirk, E. R. Reiter and A. M. Barbieri



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J.P. McGuirk, E.R. Reiter & A.M. Barbieri

Atmospheric Science Department Colorado State University Fort Collins, Co.

ABSTRACT

Hemispheric scale energy parameters have been computed and their variability on an annual time scale and a scale of a few weeks is identified and discussed.

Based on data for 15 winter seasons, the annual variability of available potential energy is linked to fluctuations of elements of the general circulation, such as the jet streams. This link establishes a relationship between simply computed hemispheric scale parameters and regional weather conditions affecting man's economic well being.

Daily statistics on three winter seasons of hemispheric energy modes have been analyzed with the ultimate goal of understanding and reproducing the behavior of the hemispheric energy cycle. Two phenomena are discussed:

- (1) A strong 22-26 day cycle in the energy modes which may be linked to synoptic scale systems and the atmospheric index cycle; and
- (2) A large midwinter "dip" in zonal available potential energy which is tentatively linked to a possible movement of the Hadley cell. This "dip" is shown to have a profound effect on the meridional temperature gradient, particularly in midlatitudes.

1. INTRODUCTION

The world food production, under the stress of an exponentially increasing world population and of vastly increased prices for energy, has become a focal point of concern. This concern was heightened by reports that indicate a deterioration of climatic conditions in the middle and high latitudes of the northern hemisphere where most of the world's agricultural activities are concentrated. All of a sudden we have come to realize that global and national economies are rather sensitive to the vagaries of weather and climate. To illustrate this point, the spectrum of kinetic energy of the zonal wind component has been plotted in Fig. 1 (after Vinnichenko and Dutton, 1969), together with crude estimates of dollar losses sustained by adverse effects of atmospheric variability. This variability is given interms of a space scale (upper part of diagram) and a time scale (lower part of diagram). Economic loss estimates of prolonged drought or flood disasters are made on a hemispheric, not national, basis; so are the estimates for a severe climatic deterioration.

Moderate climatic changes have been experienced during recorded history, especially in the middle and high latitudes of the northern hemisphere. The temperature regime in this region of the globe during the 1940's was the mildest on record within the time span of the last millennium (Fig. 2). A spectacular amelioration of climatic conditions occurred between the 1880's and the middle of the present century (see e.g., Kellogg and Schneider, 1974). This improvement in climatic conditions is also reflected in the extent of sea-ice coverage of the North Atlantic (Fig. 3).

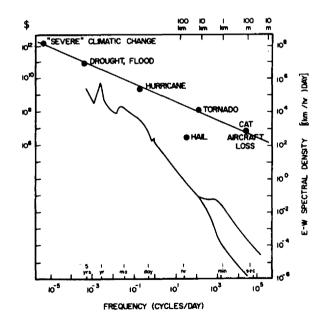


Fig. 1. Spectrum of the kinetic energy of the zonal wind component (after Vinnichenko and Dutton, 1969) units of [(km²/hr²) day] labelled on right side of diagram, and approximate dollar values of weather-related disasters (partly after Maunder, 1970), labelled on left side of diagram.

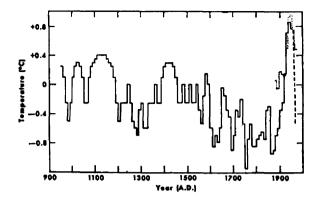


Fig. 2. Mean annual temperature in Iceland over the past millennium. (Data courtesy of P. Bergthorsson.) Dashed line shows the rate of decline in the period 1961 to 1971, dotted line indicates the variation of mean temperature in the northern hemisphere plotted to the same scale (from Bryson, 1974a).

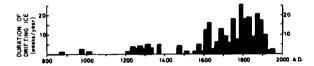


Fig. 3. Duration of drifting ice near the coasts of Iceland in units of weeks per year (after Koch, 1945, from Flohn, 1973).

It has been pointed out by Korff and Flohn (1969) that changes in the temperature of the polar regions can lead to adjustments in the meridional temperature gradient if the tropical regions do not undergo the same temperature trend. This, in turn, will lead to adjustments in the general circulation of the atmosphere, specifically in the position of the subtropical high-pressure belt. This belt will undergo a northward shift with a decreasing meridional temperature gradient (Fig. 4). We witness such a migration of the subtropical anticyclones with the progression of seasons each year from winter to summer. The climatic changes indicated in Fig. 2, of course, are much more subtle than the seasonal changes of weather patterns. Nevertheless, they are expected to leave an imprint on the behavior of the general circulation of the atmosphere.

One region of the globe where weather patterns are strongly dependent upon the seasonal migration and position of the subtropical anticyclonic belt is the Indian subcontinent. The "burst" of the summer monsoon in India, Pakistan, and Bangladesh depends on the establishment of a quasi-permanent high-pressure regime in the upper troposphere over Tibet, the shift of the westerly subtropical jet from northern India into Mongolia and China, and the establishment of an easterly jet stream over India (Koteswaram, 1958. [For a detailed discussion see Reiter, 1969b]). With a "climatic trend" by the subtropical anticyclones to undergo a poleward shift one would expect the Indian summer monsoon to burst somewhat earlier than usual; to be less susceptible to "breaks" (i.e., interruptions of the monsoonal flow patterns caused by a weakening of the anticyclone over Tibet and of the easterly flow regime over India); and to "retreat" later than normal in autumn. The

sum-total of these effects will lead to increased precipitation over the Indian subcontinent during the summer monsoon season. This is reflected in Fig. 5 which shows an abnormally low drought occurrence in India during the favorable climate period of the 1940's and 1950's.

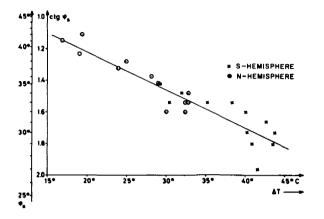


Fig. 4. Correlation between geographic ¹atitude (ϕ_S) of the subtropical high-pressure belt and the mean tropospheric (300-700 mb) temperature gradient between equator and pole (after Korff and Flohn, 1969, from Flohn, 1973).

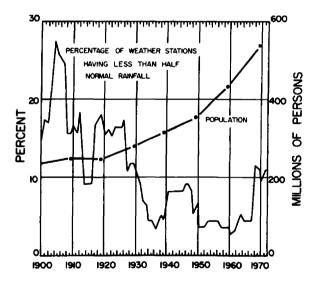


Fig. 5. Trends in the percentage of weather stations in northwestern India reporting less than half of usual rainfall in a given year. Overlapping ten-year averages (from Bryson, 1974b).

Another region of concern is the Sahel stretching along the southern rim of the Sahara desert. This region also depends on the rainfall during the summer season which originates from the moist monsoonal air currents carried inland south of the intertropical convergence zone (ITCZ) (Fig. 6). Where the ITCZ surface intersects with the earth's surface, the moist monsoonal air layer is too shallow to permit the development of convective thundershowers that penetrate into the convectively unstable, but dryer, layer aloft. Only at

some distance south of the ITCZ intersection line with the ground will there be abundant rainfall triggered by traveling disturbances in the easterly flow aloft (Mbele-Mbong, 1974).

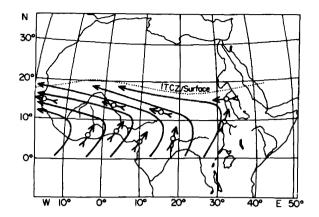


Fig. 6. Extent of monsoonal flow over Central Africa (from Bryson, 1974c).

With a "climatic" northward shift of the subtropical high-pressure belt, the African ITCZ will follow suit. The climatic optimum during the 1950's brought adequate amounts of rainfall to the Sahel. This is reflected in the water level of Lake Chad, as well as in the Nile run-off (Fig. 7).

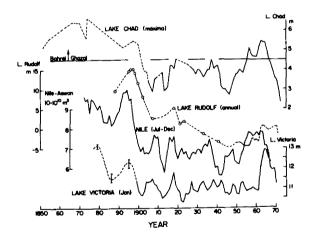


Fig. 7. Lake levels of Lake Chad, Lake Rudolf and Lake Victoria, and run-off of the Nile at Aswan (from Bryson, 1974c).

2. THE RECENT HEMISPHERIC COOLING TREND

The dense observational network in the northern hemisphere which was established in the late 1940's permits an assessment of hemispheric temperature trends. It appears that a cooling trend, especially noticeable in high latitudes, had set in around 1960 (Figs. 8 and 9) (see also Rodewald, 1972). This trend was studied in detail by Starr and Oort (1973) using the northern hemisphere upper-air data between 1958 and 1963 carefully processed at MIT.

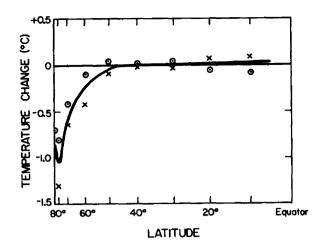


Fig. 8. Departures of the mean temperatures for the decade 1961-1970 from the mean temperatures 1931-1960 (from Flohn, 1973).

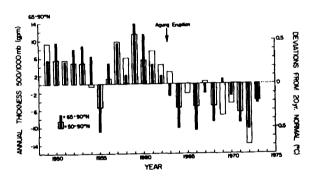


Fig. 9. Annual mean values of the thickness 500 100 mb in the latitude regions 65-900N, and 50-900N, expressed as departures from the 20 year normal values (1953-1972) (from Bryson, 1974c).

The consequences of this trend have been seen to reverberate through the precipitation regimes over India (Fig. 5) and the Sahel. The Sahelian drought, which apparently peaked in 1972, is still fresh in our minds. Unfortunately, agricultural practices and the increase in population (see also Fig. 5) adjusted to the climatic optimum of the 1940's and 1950's. Re adjustment to a somewhat harsher climate will bring grave economic, social and political stresses, not only to the areas immediately affected by adverse precipitation trends but also to the world at large which has not yet learned how to stabilize its demand for non-renewable resources.

Alarmed by the prospects of a deteriorating climate, and by the suggestions made by several scientists that this deterioration originated from such anthropogenic causes as air pollution by suspended particulates, we decided to check the northern hemisphere temperature trend with our NMC data bank. These data are less carefully screened than the "MIT period" of five years, but they offer the advantage of a longer time record which, eventually, will be brought up to current.* (Unfortun-

*We are indebted to NCAR for supplying us with copies of the NMC tapes. Dr.R.F. Adler, now with NASA Goddard Space Flight Center, performed the original analyses of the hemispheric temperature trends. Mrs. Alice Fields was in charge of programming and computer processing of the data.

ately the data tapes for 1961 and 1962 could not be obtained. Thus, a gap exists in our analyses for these two years.)

Fig. 10 reveals temperature variations at 500 mb for the winter months and for the "cold season" which are quite similar to those shown in Fig. 9. Especially the "precipitous" decline between 1959 and 1963, the slight recovery between 1964 and 1968, and the further decline until 1972 are well reflected.

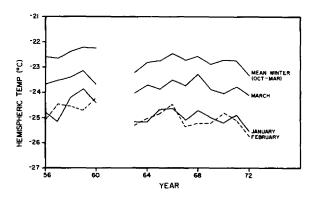


Fig. 10. Mean hemispheric temperatures north of 200N at $500~\mathrm{mb}$.

From Fig. 11 we can see that the mid-winter cooling trend was strongest at high latitudes, especially between 1959 and 1964. The mid-latitude belt between 30°N and 60°N also reveals a slight irregular cooling trend between 1959 and 1972 for January and for the cold season, less so for March (Fig. 12). This trend is less pronounced here than at high latitudes (Fig. 11). From this we should expect a gradual increase of the zonal available potential energy in the northern hemisphere and, as a consequence, a strengthening of the westerly circulation in the jet stream belt of middle latitudes, together with a slight southward shift of this belt.

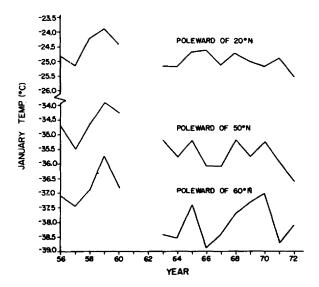


Fig. 11. Mean January temperatures for various spherical segments of the northern hemisphere.

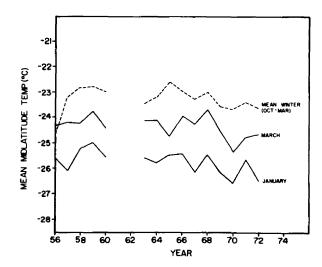


Fig. 12. Mean temperatures for latitude belt $30^{\circ}N$ to $60^{\circ}N$ for time periods as indicated.

The zonal available potential energy, averaged over the month of January, underwent large interannual fluctuations within the time period 1956-1972 as shown in Fig. 13. Discounting these fluctuations, an increasing trend is noticeable, in agreement with our hypothesis stated earlier. The zonal kinetic energy, on the other hand, does not reveal a significant longterm trend, although analysis of kinetic energy only at 500 mb does not allow us to say with any certitude that a long-term trend does not exist. The eddy-available potential and kinetic energies (Fig. 14) reveal no significant trends either, except that during the two years of missing data (1961 and 1962) a significant increase in both energies seemed to have occurred. The cooling trend which was relatively steep during this short time period (Fig. 11) might have led to an adjustment in the general circulation resulting in an increased planetary wave activity.

In order to check the hypothesis that the cooling trend in high latitudes might have resulted in a southward shift of the jet stream belt, we computed the "cold-pool area" bounded on its equatorward side by the position of the polar front. After an inspection of a large number of daily 500-mb maps we decided that the location of the -30°C isotherm at this pressure level adequately defines the mean position of the polar front during the cold season (see also Palmén and Newton, 1969; Reiter and Macdonald, 1973; Reiter, 1975).

The upper part of Fig. 15 shows the results of our computations for January. Indeed, we find a somewhat irregular trend between 1959 and 1972 that suggests a gradual increase of the area inside the -30°C isotherm at 500 mb. This can be interpreted as a slight equatorward shift of the polar front and its associated jet stream.

The area-weighed mean temperature inside the -30°C isotherm at 500 mb defines the "cold-pool temperature." This parameter, computed for the month of January, is shown in the lower part of Fig. 15. Two facts become apparent: There exists a reasonably good correlation between the cold-pool temperature and the hemispheric mean temperature poleward of 50°N (middle curve in Fig. 11, see also Fig. 16). Secondly, the sporadic interannual fluctuations of the cold-pool temperature and the cold-pool area (lower and upper curves in Fig. 15)

are poorly correlated with each other. Thus, there are winters -- or, more properly, Januaries -- when the cold pool is large but relatively warm, and other winters when it is relatively cold.

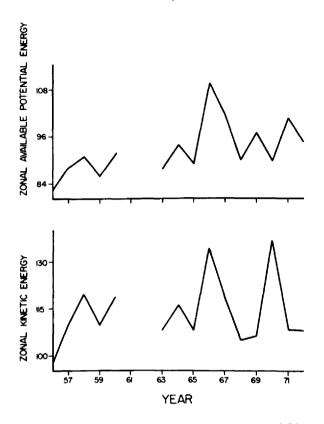


Fig. 13. Mean zonal available potential and kinetic energies (arbitrary units) for January, computed for the northern hemisphere poleward of 20°N and for the 500-mb layer.

The correlation between mean cold-pool temperature in January and mean zonal available potential energy scatters considerably. All data points for the period under investigation lie to the right of line shown in Fig. 17, which appears to set a limiting condition between the cooling of the polar regions and the buildup of hemispheric potential energy.

In Fig. 18 we have plotted the mean meridional temperature profiles at 500 mb for the months of January 1956 to 1972 (1961 and 1962 are missing). It is easily seen that the largest interannual variability is confined to high latitudes, 1963 showing the coldest conditions there. In middle latitudes, on the other hand, January of 1972 established the cold record for the period under investigation. During this year the entire meridional temperature gradient was confined to latitudes south of 65°N. We should expect, therefore,

that the jet-stream activity during this winter month also had shifted relatively far to the south. We should recall that this year was characterized by a severe drought in the Sahel. Drought conditions also were wide spread in India.

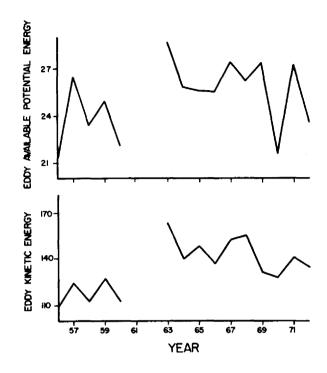


Fig. 14. Mean eddy available potential and kinetic energies (arbitrary units) for January, computed for the northern hemisphere poleward of 200N and the 500-mb layer.

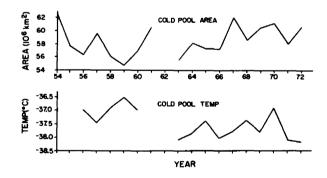


Fig. 15. Area of the cold pool (10^6km^2) bounded by the -30°C isotherm at 500 mb, and mean areaweighted temperature of this area at 500 mb.

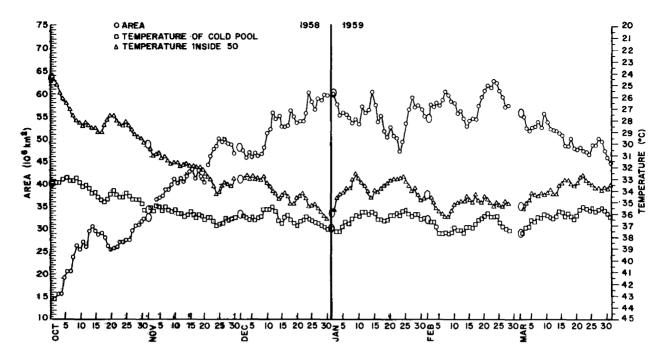


Fig. 16. Daily values for the 1958-1959 cold season of the cold-pool area at 500 mb (i.e., the area bounded by the -30° C isotherm at 500 mb), the area-weighted temperature of this area at 500 mb, and the mean temperature of the spherical segment north of 50° N.

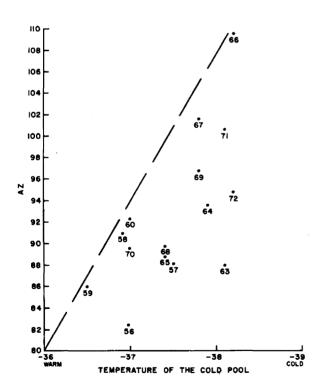


Fig. 17. Mean cold-pool temperature (i.e., mean temperature of the area bounded by the -30°C isotherm at 500 mb) and zonal available potential energy (arbitrary units) for years indicated by two digits.

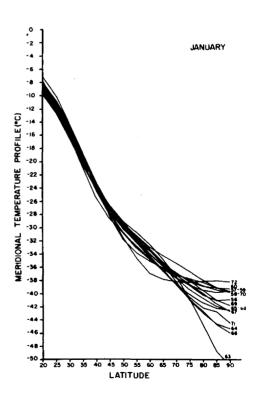


Fig. 18. Mean meridional temperature profiles at 500 mb for years indicated by last two digits.

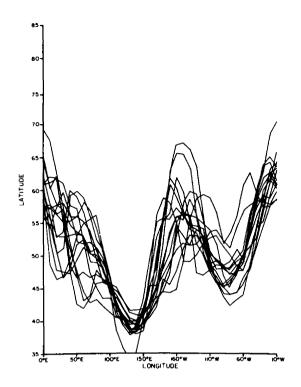


Fig. 19. Average latitudinal position of the -30°C isotherm at 500 mb during January as a function of longitude, for fifteen years 1956-1972 (1961 and 1962 are missing).

3. REGIONAL EFFECTS

Starr and Oort (1973) and Landsberg and Albert (1975) showed that the temperature trend experienced in the northern hemisphere during recent history did not spread uniformly around the globe. The warming that took place between the latter part of the 19th century and the "normal period" of 1941-1970 was confined in the North American sector to the U. S. Rocky Mountain States and to the Midwestern "cornbelt" between Nebraska and Indiana. Surface temperature rises of over 2°C were observed in these regions, whereas most of the southwestern United States and the Pacific coastal regions remained unaffected by this warming trend.

Landsberg and Albert (1975) furthermore showed that a change in annual mean temperatures can have a profound effect on the length of the freeze-free growing season, especially in regions with strong topographic structuring. In Oregon, for instance, a two-degree rise in the annual mean temperature would result in an average increase of more than 50 days in the length of the growing season. In North Dakota, on the other hand, a similar change in mean temperature would produce a change in the growing season of perhaps only 10 days.

If the long-term warming and cooling trends reveal marked regional preferences, we have to suspect the same for the interannual variability of the "cold pool" of the northern hemisphere. This regional behavior can be expressed by the variation in latitude at which the -30°C isotherm at 500 mb occurs at different longitudes. Average latitude positions of this isotherm during January are shown in Fig. 19. The standard deviation in the latitudinal fluctuations of this isotherm for the 15-year period 1956-1972 and for the average of January, is shown in Fig. 20 (1961 and 1962 missing).

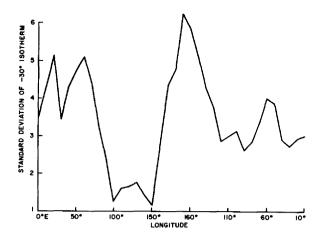


Fig. 20. Standard deviation of the data sample shown in Fig. 19.

Minimum variability is found between 100°E and 150°E, i.e., in the Mongolian and Japanese longitude sector. Here, apparently, minor climatic trends in the hemispheric winter-time temperature regimes will have little regional impact.

Maximum variability is found over the Pacific Ocean in the Guld of Alaska at 160°W. We have to suspect that the interannual variability of the mean January position of the polar front in this region, together with regional sea-surface temperature anomalies (see e.g., Namias, 1969) might produce an interannually variable forcing of the general circulation, albeit small, through the sensible and latent heat input into the atmosphere. This forcing, if it is present to a significant degree, is expected to be of a rather complex nature and should be most effective in the lower troposphere. An attempt to correlate the meridional heat flux at 60°N and at 500 mb in this region of the Pacific with the latitudinal position of the -30°C isotherm at 170°W reveals a tendency for decreased flux as the polar front moves to the south of this parallel circle (Fig. 21). The data points for January 1958 and 1967 disturb this trend. During these months the -30°C isotherm was found exceptionally far to the south, indicating the presence of strongly anomalous synoptic conditions in the Gulf of Alaska.

A correlation attempt as shown in Fig. 21, in retrospect, is not the best in considering regional forcing effects on the atmospheric circulation. If seasurface temperature anomalies, together with atmospheric circulation anomalies, were indeed producing a noticeable effect on atmospheric energetics or on flow patterns, we should expect to find such effects more pronounced downstream, perhaps over North America, than in the same longitude sector. This aspect remains to be investigated. At this point of our investigation we can only stress the large interannual temperature variability in the Gulf of Alaska during January.

From Fig. 20 it also appears that large interannual variations of the -30°C isotherm at 500 mb and during January exist near 60°E and 60°W. These longitudes intersect Kasakhstan and the Labrador coast, respectively. The agriculturally productive sector of the United States between 80°W and 100°W shows less interannual variability in the January temperature regime at 500 mb than does the European sector, including the Ural Mountain region. We have not yet studied the

possible effects of winter weather on agricultural crops, such as winter wheat, nor have we completed our analyses for the months characteristic of the planting, growing and harvesting seasons. Work along these lines is in progress, however. Results, if they are similar to those presented in Figs. 19 and 20, might help in identifying those regions of the northern hemisphere where agricultural productivity might be most vulnerable to the interannual variability of weather regimes and to climatic trends.

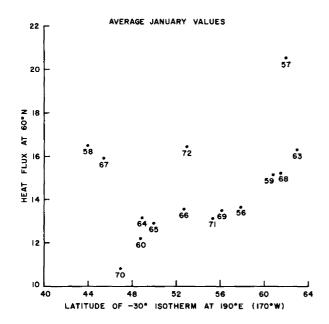


Fig. 21. Heat flux across $60^{\circ}N$ and the latitudinal position of the -30°C isotherm at 500 mb and at 170°W for January of the years indicated by last two digits.

We have to realize, of course, that the temperature regimes at the 500-mb surface cannot be readily translated into an equivalent behavior of surface temperatures. Additional factors, such as cloudiness, snow cover, etc., exercise their influence on the heat budget, hence the temperatures of the lowest layers of the troposphere. A certain correlation between the 500-mb temperature regime and surface temperatures does exist however. From Fig. 22 we can see that the Moscow region shows a correlation between the latitude of the -30°C isotherm at 500 mb during January and the mean surface temperature of January, although the data points scatter considerably. Pittsburgh (Fig. 23) shows less scatter, especially for those years in which the meridional temperature gradient over the southern U. S. was relatively tight. The term "tight" is arbitrary in that the data were stratified into two categories: The eight years with the steepest meridional temperature gradient between 25°N and 30°N which were classified as "tight"; and the remaining seven years which were classified "loose."

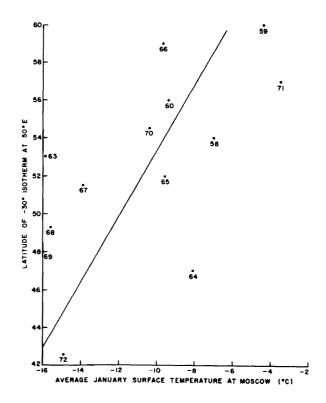


Fig. 22. Mean January surface temperature at Moscow, USSR, and the latitude of the -30°C isotherm at 500 mb and at 50° E for years indicated by last two digits.

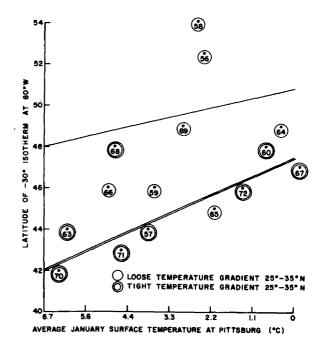


Fig. 23. Mean January surface temperature at Pittsburgh, Pennsylvania and the latitude of the -30° C isotherm at 500 mb and at 80° W for years indicated by last two digits.

4. HEMISPHERIC SHORT TERM VARIABILITY

The zonal and eddy modes of the hemispherically averaged available potential and kinetic energies undergo rather large fluctuations during a single season. It is desirable to investigate these fluctuations for two reasons:

- (i) The understanding of the interactions of these energy modes within a season will be a powerful diagnostic tool in understanding long-range weather fluctuations.
- (ii) Understanding the interannual variability in the seasonal behavior of these energy modes will bring us a step closer to the successful modeling of climatic trends.

Towards these ends we computed the four modes of energy for the cold seasons from NMC tapes. The energy parameters, in units of 10^5 joules/m², were computed for the spherical segment north of $20^\circ N$, and for the layer 1000 mb to 100 mb. Three winter seasons, 1968 to 1971, have been analyzed so far. Work is in progress to extend this analysis for the period 1963 to present. Typical data for zonal and eddy available potential energy, and for zonal and eddy kinetic energy are shown in Figs. 24 to 35. Three-day and thirty-one-day, unweighted running averages were used as a simple highpass and low-pass filtering technique. We avoided using monthly average values because they may obscure significant short-term fluctuations in the data.

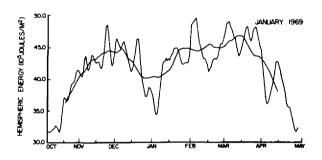
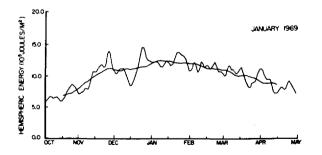


Fig. 24. 3-day and 31-day running means of zonal available potential energy (units 10^5 joules/m²) for the period October 1968 through April 1969 and for the region north of $20^{\rm O}N$ and the layer 1000 to 100 mb.



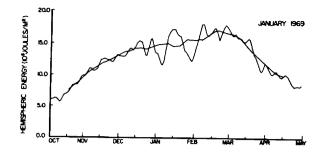


Fig. 26. Same as Fig. 24, except for zonal kinetic energy.

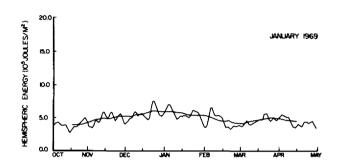


Fig. 27. Same as Fig. 24, except for eddy kinetic energy.

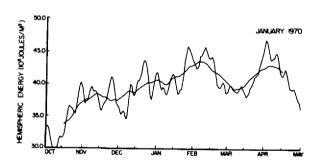


Fig. 28. Same as Fig. 24, except for period October 1969 through April 1970, zonal available potential energy.

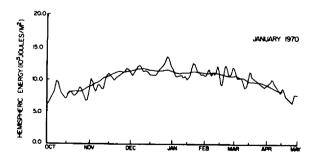


Fig. 29. Same as Fig. 28, except for eddy-available potential energy.

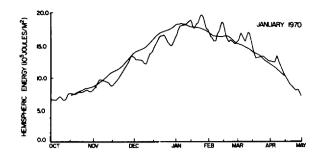


Fig. 30. Same as Fig. 28, except for zonal kinetic energy.

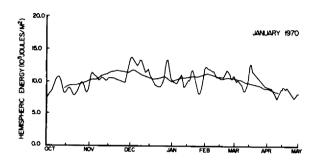


Fig. 31. Same as Fig. 28, except for eddy kinetic energy.

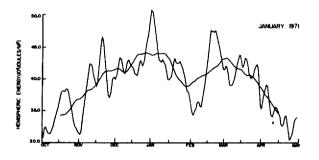


Fig. 32. Same as Fig. 24, except for period October 1970 through April 1971, zonal available potential energy.

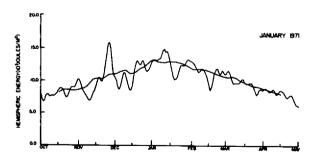


Fig. 33. Same as Fig. 32, except for eddy-available potential energy.

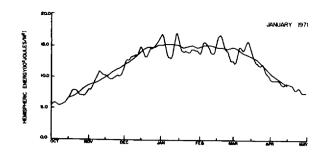


Fig. 34. Same as Fig. 32, except for zonal kinetic energy.

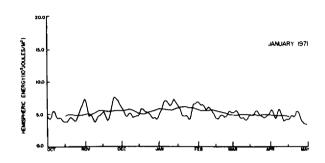


Fig. 35. Same as Fig. 32, except for eddy kinetic energy.

The most obvious characteristic of the data is the large "dip" in AZ in midwinter, unaccompanied by any major change in AE, KZ, KE.* Such a dip occurs in all three years analyzed so far with the relative minima appearing 27 December 1968, 7 March 1970, and 20 January 1971. This phenomenon has been reported previously (Krueger et al, 1965), but very recently has been called into question (Peixoto and Oort, 1974). This conflict in opinions, as well as some possible explanations, are discussed later in some depth.

The second interesting feature of the data is the rather large short-term variation of individual hemispheric energy components. Particularly, KE and AE undergo changes of nearly 100% in just 8 or 10 days. Masked by its comparatively larger magnitude, AZ often undergoes changes greater in magnitude than the eddy potential or kinetic energy in less than a week.

These comparatively large-amplitude fluctuations in the energy modes reveal significant phase relationships among the four parameters. This is evident in Fig. 36-38 which show typical three-day average deviations from 31-day averages. Of particular interest are the following correlations:

^{*}The following notation has been adopted (Reiter, 1969a).

AZ = Zonal Available Potential Energy

AE = Eddy Available Potential Energy

KZ = Zonal Kinetic Energy

KE = Eddy Kinetic Energy

- 1) AZ and KE appear nearly out of phase;
- AZ and AE bear an even stronger out-of-phase relationship;
- 3) KE and AE are very close in phase to each other with KE leading AE by one or two days, contrary to expectations from conventional energy flow diagrams (see e.g., Reiter, 1969a for references and examples) which show AE feeding KE.

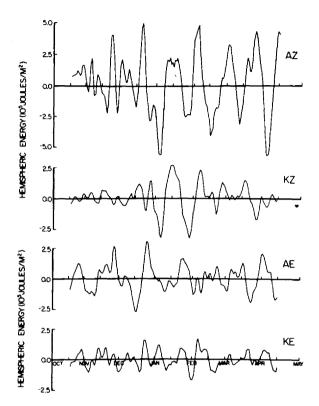


Fig. 36. Deviations of the 3-day average energy modes from their 31-day averages for the period October 1968 through April 1969.

Quasi-periodicities in the fluctuations of the four modes of energy, and correlations between them can be examined readily from auto- and cospectra of these energy parameters. An example is shown in Fig. 39. The points of each curve represents the contribution to the total variance by that particular harmonic frequency. The individual curves represent the average of the spectra of the three seasons October through April of 1968 to 1971. These individual curves were constructed from appropriately smoothed Fourier transforms of the correlation functions.

The spectral power peak at about 24 days exceeds the 90% confidence level computed by the chi-square technique. The record length was too short to adequately define or establish the significance of the seasonal harmonic at about 90 days. In order to arrive at a significant peak in this frequency domain, one would have to include the summer months into the analysis. This might, however, alter the peak at 24 days, if during summer the characteristic energy cycles run on a somewhat different frequency by virtue of the reduced temperature gradient between equator and pole. The subharmonics of the 22-24 day cycle at 12 or 13 days and at seven or eight days, while present, do not seem highly significant. The spectrum for KZ was not extended below a period of 54 days because no peaks above

the noise level were present. The implication is that KZ follows only the gross behavior of AZ and the total zonal kinetic energy is not too sensitive to tropospheric disturbances.

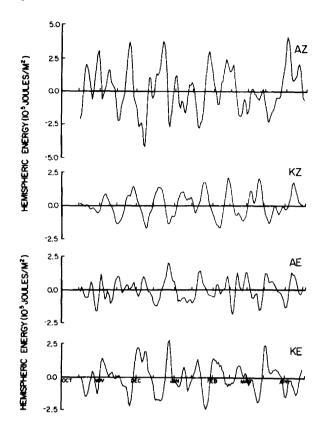


Fig. 37. Same as Fig. 36, except for time period October 1969 through April 1970.

The relationship between AZ and KE becomes quite evident under cross-spectrum analysis. In Fig. 40, the coherence, a measure of the correlated amplitudes of both signals, AZ and KE, at a given frequency, and the spectral gain (cospectrum) which represents the amplitude of the cospectrum, are plotted. In mathematical terms the coherence is the ratio of the gain and the square root of the product of the individual autospectral power densities. At a given frequency if both signals (AZ and KE) are perfect sine waves, the coherence is equal to one; if both signals are uncorrelated white noise, the coherence should approach zero. The gain is the root-mean square of the real and imaginary parts of the cross spectrums of AZ and KE. If the gain is significantly large at a given frequency, either AZ or KE, or both AZ and KE have a strong harmonic at that frequency. If, in addition, the coherence is large (close to one), the harmonics of AZ and KE at that frequency have a strong linear dependence on each other. This also implies that while the two signals may differ in phase, this phase difference is approximately constant. Fig. 40 shows large values of gain at a frequency corresponding to a 25-day oscillation. This means that both AZ and KE vary significantly at this The coherence at the 25-day oscillation is also large (0.9) so that the variation of AZ and KE are strongly dependent on each other. The amplitude of the cospectrum (or gain) is almost as large as the seasonal amplitude. Also relatively strong are the subharmonics at 12 or 13 days. Fig. 38 shows that AZ leads KE by about 11 days, implying that AZ and KE are almost out

of phase with respect to the 25-day vacillation. This is to be expected from the form of the conversion term <AZ, AE> (see Reiter, 1969). The cospectra for the other two years show similar but slightly smaller peaks at about 25 days. We are presently investigating if this rather significant vacillation reflects itself in the "index cycle" of the northern hemisphere (see e.g., Quinet, 1974).

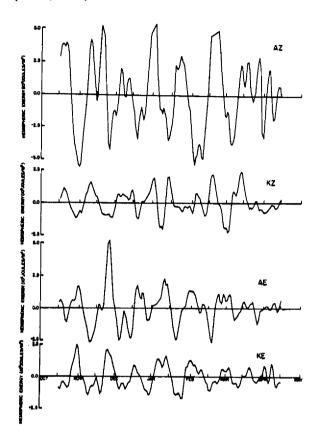


Fig. 38. Same as Fig. 36, except for time period October 1970 through April 1971.

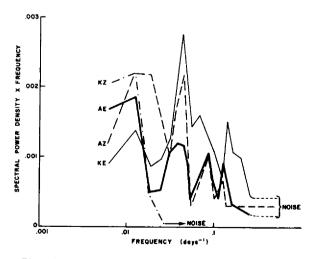


Fig. 39. Mean spectral power density of energy modes averaged over three winter seasons (October through April) from 1968 to 1971.

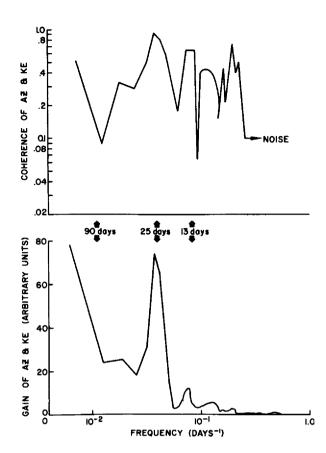


Fig. 40. Coherence and spectral gain of zonal available potential energy (AZ) and eddy kinetic energy (KE) for the winter season October 1968 to April of 1971.

Because of this strong phase relationship between AZ and KE, it seems evident that synoptic scale events (as measured by fluctuations in KE) can be directly related to the zonally averaged available potential energy. This is also evident from the short-term variations of mean meridional temperature profiles (Fig. 41).

These profiles represent the average temperature of the layer 1000 to 100 mb. The curve for 10 February 1970 is representative for conditions prior to the decrease in AZ (Fig. 28). As the zonal available potential energy decreases, the middle latitudes reveal a marked cooling trend and high latitudes exhibit a warming. We have not yet ascertained in detail to what extent the development of blocking patterns in the troposphere plays a role in this AZ "dip." A cursory examination of weather maps for the three seasons analyzed did not provide conclusive evidence for or against the significance of blocking.

A close relationship has been found between the magnitude of KE and the eddy-heat flux at some representative latitude (data will be supplied in a subsequent report).

Fig. 38 also reveals a relationship between AE and KE. A close correlation between these two energy modes should be expected because both are caused by the same physical process: the growth of unstable baroclinic

waves. The interesting phenomenon, however, is that KE leads AE. This is contrary to typical atmospheric energy-flow diagrams which show AZ being converted to AE, which in turn is converted to KE. Thus at any time, the magnitude of AE seems to be governed by the conversion <AZ, AE> which, in turn, is strongly controlled by the magnitude of KE (Fig. 42). In a developing wave, AE seems to be converted very rapidly to KE while the frictional dissipation acts more slowly.

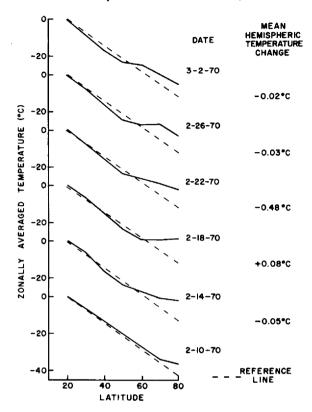


Fig. 41. Mean meridional temperature profile of the layer 1000 to 100 mb, and mean temperature change of spherical section north of $20^{\circ}N$ for selected days surrounding a typical mid-winter decrease in AZ. The dashed line provides an arbitrary constant temperature profile.

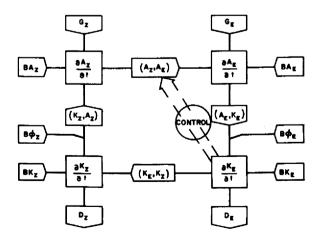


Fig. 42. Atmospheric energy transformation (see Reiter, 1969). There seems to be a feedback mechanism acting between KE and the transformation <AZ, AE>.

5. THE MIDWINTER DECREASE IN ZONAL AVAILABLE POTENTIAL ENERGY

As mentioned previously, a large drop (10-20%) is noted in AZ during mid to late winter. Such a drop has recently been questioned by Peixoto and Oort (1974). The present data reveal a drop in AZ rather conclusively. Kruger et al. (1965) also find a significant midwinter dip in AZ. Besides the difference in time periods, the dissimilarities between the present data and those cited are:

- 1) We use 3- and 31-day running averages, whereas Peixoto and Oort use fixed monthly values averaged over five years.
- We had to terminate our analyses at 20°N whereas Peixoto and Oort could extend their data base to 10°S.

It is questionable whether cumulative monthly averages will depict with any clear definition an event whose periodicity is one and one-half months or less. As Peixoto and Oort readily admit, the difference between their results and those obtained by Kruger et al. (1965) may be explained in part by the boundary flux at There is some evidence that this is indeed the We can hypothesize that the total energy flux across 20°N is modulated by a latitudinal movement of the Hadley cell. If this were true one would anticipate similar changes in the flux at 10°S, which should appear in Peixoto and Oort's data but are not evident there. If the large potential energy flux of the direct meridional Hadley cell decreases at latitude 20°N, the eddy heat transport in midlatitudes will rapidly relax the N-S temperature gradient.

Because the definition of a reference state of mean temperature as a function of height enters into the computation of available potential energy (Dutton and Johnson, 1973) AZ is particularly sensitive to small movements of the Hadley cell. For example, small netmass fluxes northward across latitude 20°N will increase AZ simply by increasing the mean surface pressure. But even larger increases occur because of the relatively warmer air and the mean lower stability south of 20°N. To estimate the magnitude of this effect, data for February from Oort and Rasmussen (1969) were used to calculate the change in AZ north of 20°N that would result by including mean temperatures as a function of height in the latitude belt between equator and 20°N, and thus arrive at a new reference state of potential energy in the northern hemisphere. It was determined that AZ of the northern hemisphere increases 4% due to the different mean temperature and 2.5% due to the different static stability that resulted from an inclusion of the latitude belt 0° to 20°N. In other words, the mass moving northward across 20°N will possess approximately 6.5% more AZ simply by virtue of possessing a different reference state. Since changes in the mean meridional circulation are not detectable by the available NMC data, indirect evidence is presented to support the hypothesis of meridional flux across 20°N.

Sample calculations were performed for the period ll to 24 February 1970, in which AZ decreased by 17% and the total available energy (AZ + AE + KZ + KE) decreased by 9.5 x 10^5 joules/m². If no energy flux across $20^\circ N$ existed, this energy decrease would have to appear as dissipation into heat and, as a consequence, the total potential energy $f_{\text{PC}_{\text{D}}}$ Td (Volume) would have to increase. However, during this period, the total potential energy also decreased by two watts/m². We, thus, are faced with an "unexplained" energy decrease of nearly three watts/m². In addition to this, Oort

and Rasmussen estimate the steady-state frictional dissipation to be 2.3 watts/ m^2 . Therefore a decrease of somewhere between three and five watts/ m^2 must be explained. Still ignoring the boundary flux across 20°N, two possible explanations exist:

- The steady-state radiation balance can be interrupted by large albedo changes or by increases of long-wave radiation;
- 2) Energy can be absorbed in evaporation of snow and ice.

We have ignored surface heating effects because they are deemed to act slower than the observed decreases in total atmospheric energy. Cursory analyses of satellite photographs do not indicate any marked changes of snow or cloud cover during the periods of energy decrease. (See U. S. Dept. of Commerce ESSA 9 data, 1971.) Therefore, it seems the "loss" in total energy cannot be accounted for by albedo changes. Furthermore, typical hemispheric values of the generation of (AZ + AE) are about six watts/m². The "unaccounted" energy decrease is almost as high, namely three to five watts/m². Albedo would have to vary by 50 to 100% if it were held responsible for this loss in energy due to the energy-generation terms (Reiter, 1969). Such variations, if indeed present, would be easily detected from satellite photographs.

Using areal extent of snow cover (Kukla and Kukla, 1974) an estimated evaporation of about one cm of water over this whole area would account for the observed energy deficit. If we assume, however, that evaporation takes place significantly only in the area over which the 0°C sea-level isotherm retracts toward the north, a water column of 20 cm would have to evaporate to make up for the energy loss.

Using Oort and Rasmussen's data, the following estimates for fluxes at $20^{\circ}N$ in February can be made (units are in terms of watts per m^2 of the total area north of $20^{\circ}N$):

Sensible heat : 20 w/m^2 southward Potential energy : 46 w/m^2 northward Latent heat : 32 w/m^2 southward Net flux : 6 w/m^2 southward

Thus it appears that net energy fluxes southward occur typically in late winter and possess magnitudes similar to those of the "unexplained" decreases in AZ.

If there is a substantial southward net flux of energy at 20°N we should expect a cooling in the area north of 20°N. Fig. 41 shows such a net cooling between 18 and 22 February 1970. From 10 February to 18 February 1970, the hemispheric temperature changes by about +0.03°C but a fairly dramatic redistribution of hemispheric temperatures has occurred. The midlatitudes from 35°-60°N have cooled while the pole has warmed markedly. Daily synoptic maps confirm this redistribution by showing a strong surge of tropical air flowing across the pole from a degenerating Ω -block. From 18 February through 2 March the entire hemisphere cools. In February 1970 the large decrease in AZ was accompanied by temperature drops of nearly 4°C in midlatitudes.

6. Modeling of the Energy Cycle

One of the primary reasons for investigating the spectral properties of the energy cycle is to gain sufficient insight into the interactions between the various energy components to allow some form of numerical simulation modeling. The modeling philosophy is based on two hypotheses:

(i) The atmosphere functions like a "dynamic system" in which the rapid increase of potential energy in late fall triggers an "overload-release mechanism" which drains the available potential energy. This mechanism is offered by unstable barcclinic waves. In the atmosphere it is highly over efficient in transporting heat poleward. It takes a certain time after such barcclinic energy releases before the hemispheric potential energy supply can build up again to a quantity sufficient to trigger the release mechanism.

(ii) This "dynamic system" can be empirically approximated by a system of forced-damped oscillators of the form:

$$\frac{dN}{dt} = -\lambda N + a + \Sigma b_{i} \sin c_{i} t \qquad (1)$$

where N is an energy mode; a, b_i , c_i are empirical forcing terms. The first hypothesis is not new, having been suggested in slightly different forms by Kraus and Lorenz (1966) and Namias (1950). Its application to reproduce atmospheric general circulation variations should prove to be of interest.

The second hypothesis raises two questions:

- (1) Can the energy cycle be modeled by a simple oscillator? By taking the output of the electrical analog of Kraus and Lorenz (1966), which looks very similar to a smoothed time series of AZ, it is possible to construct a differential equation similar to Eq. (1) which will approximate our observational data. Therefore, it seems quite probable that equations of the form Eq. (1) will be adequate for a simple model of the general circulation.
- (2) Will the set of equations that fits our opservational data be unique? The simple answer is no. As in any empirical model the assumption will be made that if external forcing to the atmosphere does not change greatly, two different sets of equations describing the same data set should produce approximately the same results.

This model is presently in its early design $phase\cdot$ It incorporates the following features:

- (i) Four inhomogeneous, coupled, first-order, nonlinear equations describing the time-change of the four modes of energy.
- (ii) Because of the analytic form of the conversion term from AZ to AE, which depends on the meridional eddy velocity, it is hypothesized that this conversion will be controlled by the amplitude of the eddy kinetic energy. Thus, as suggested in Fig. 42, one or more positive feedback mechanisms will be incorporated into this model.

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Authors: J.P. McGuirk, E.R. Reiter and A.M. Barbieri

551.513.1

Authors: J.P. McGuirk, E.R. Reiter and A.M. Barbieri

551.513.1

ON THE VARIABILITY OF HEMISPHERIC SCALE ENERGY PARAMETERS

Colorado State University
Department of Atmospheric Science

Subject Headings: General Circulation Potential Energy Kinetic Energy Polar Front Subject Headings: General Circulation Potential Energy Kinetic Energy Polar Front

Environmental Research Paper No. . 1, January 1975. 15p.

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No. AT(11-1)-1340

Hemispheric scale energy parameters have been computed and their variability on an annual time scale and a scale of a few weeks is

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Colorado State University
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Subject Headings: General Circulation Potential Energy Kinetic Energy Polar Front

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